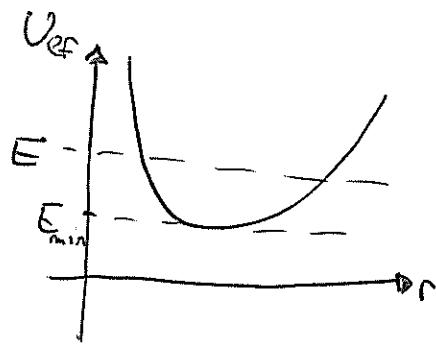


## Auxiliar 22

P1]  $U_{\text{ef}} = \frac{\ell^2}{2mr^2} + Cr^3$



a)  $F(r)$ ?

b) Radio  $r_0$  de órbita circular

c) Cuando  $E > E_{\min}$ , encontrar  $\frac{dr}{dt}$  en  $r=r_0$ .

Expresar en  $E, E_{\min}, \ell, C_m, r_0$

d)  $K$  en  $r=r_{\max}$

Sol: a) El potencial:  $U(r) = Cr^3$

$$\Rightarrow F(r) = -\frac{d}{dr}U \Rightarrow \boxed{F(r) = -3Cr^2} \quad \left. \begin{array}{l} \\ \end{array} \right\} \text{atractiva}$$

b) Para que sea órbita circular:

$$\frac{dU_{\text{ef}}}{dr} \Big|_{r_0} = 0 = \frac{-\ell^2}{mr_0^3} + 3Cr_0^2 \rightarrow 0 = -\ell^2 + 3Cmr_0^5$$

$$\Rightarrow \boxed{r_0 = \left( \frac{\ell^2}{3Cm} \right)^{1/5}}$$

c) La energía cinética:

$$K = \frac{\ell^2}{2mr^2} + \frac{m}{2} \left( \frac{dr}{dt} \right)^2$$

$$\text{en } r_0 \rightarrow E = \frac{\ell^2}{2mr_0^2} + \frac{m}{2} \left( \frac{dr}{dt} \right)^2 \Big|_{r_0} + U_{\text{ef}}(r_0)$$

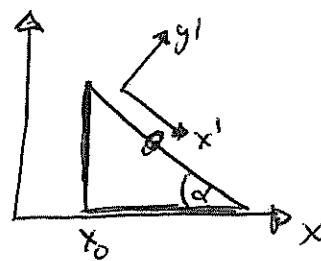
$$E_{\min} = \frac{\ell^2}{2mr_0^2} + U_{\text{ef}}(r_0)$$

$$\Rightarrow E - E_{\min} = \frac{m}{2} \left( \frac{dr}{dt} \right)^2 \Big|_{r_0} \Rightarrow \boxed{\left| \frac{dr}{dt} \right| = \sqrt{\frac{2}{m}(E - E_{\min})}}$$

d) En  $r_{\max}$ ,  $\left. \frac{dr}{dt} \right|_{r_{\max}} = 0$

$$\Rightarrow \boxed{K = \frac{\ell^2}{2mr_{\max}^2}}$$

P2)



En  $t=0$  el anillo está en reposo.

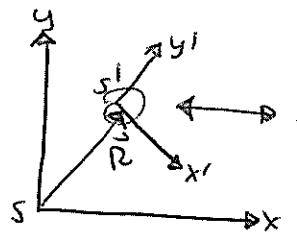
La estructura oscila con:

$$x(t) = A(1 - \cos(\omega t))$$

a) Ec. en coor.  $x'$       b) Normal en  $y'$

c)  $x'(t)$  con  $x'(t=0)=0$

Sol:

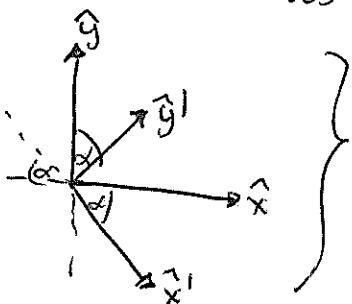


$$\left. \begin{array}{l} \vec{R} = y_0 \hat{y} + A(1 - \cos(\omega t)) \hat{x} \\ \Rightarrow \vec{R} = A \sin(\omega t) \hat{y} \\ \vec{R} = A \omega^2 \cos(\omega t) \hat{x} \end{array} \right\} \vec{R} = 0$$

En el sist. no inercial:

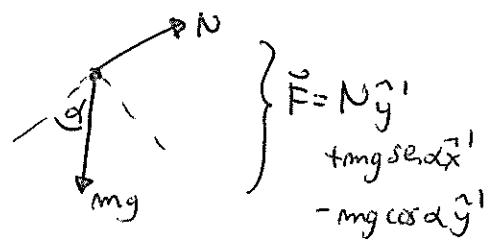
$$\vec{x}' = x' \hat{x}', \quad \vec{v}' = \dot{x}' \hat{x}', \quad \vec{a}' = \ddot{x}' \hat{x}'$$

Relaciona los vectores unitarios:



$$\left. \begin{array}{l} \hat{x} = \cos \alpha \hat{x}' + \sin \alpha \hat{y}' \\ \hat{y} = -\sin \alpha \hat{x}' + \cos \alpha \hat{y}' \end{array} \right\}$$

DCL



$$\Rightarrow m \ddot{x}' = mg (\sin \alpha \hat{x}' - \cos \alpha \hat{y}') + N \hat{y}' - Am \omega^2 \cos(\omega t) (\cos \alpha \hat{x}' + \sin \alpha \hat{y}')$$

en  $\hat{x}'$ :

$$\boxed{\ddot{x}' = g \sin \alpha - A \omega^2 \cos(\omega t) \cos \alpha} \quad (1)$$

b) en  $\hat{y}'$ :

$$0 = -mg \cos \alpha + N - A m \omega^2 \cos(\omega t) \sin \alpha$$

$$\Rightarrow \boxed{N = m (g \cos \alpha + A \omega^2 \cos(\omega t) \sin \alpha)}$$

c) Resolvemos (1):

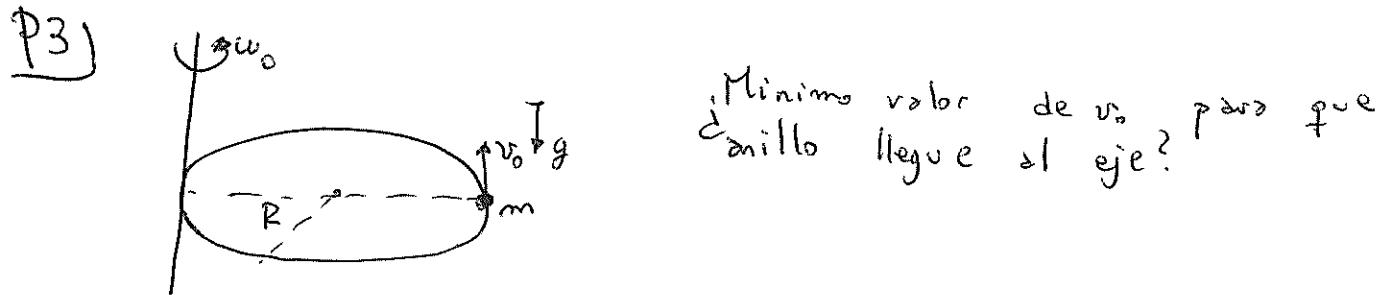
$$\ddot{x}' = \frac{d \dot{x}'}{dt}$$

$$\int_0^t d \dot{x}' = \int_0^t (g \sin \alpha - A \omega^2 \cos(\omega t) \cos \alpha) dt$$

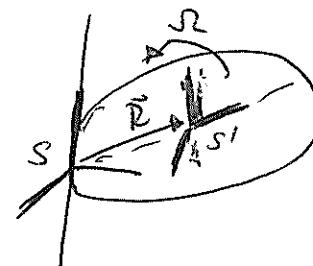
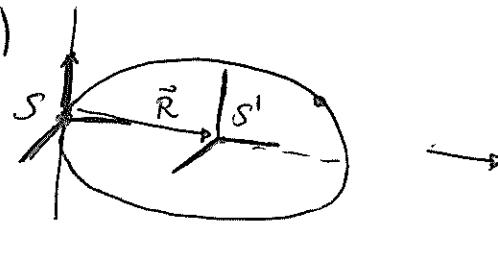
$$\dot{x}' = g \sin \alpha t - A \omega \cos \alpha \sin(\omega t) \Big|_0^t = g \sin \alpha t - A \omega \cos \alpha \sin(\omega t)$$

$$\rightarrow \ddot{x}' = \frac{d \dot{x}'}{dt} \Rightarrow \int_0^t d \dot{x}' = \int_0^t (g \sin \alpha t - A \omega \cos \alpha \sin(\omega t)) dt$$

$$\Rightarrow \boxed{x'(t) = \frac{g}{2} \sin \alpha t^2 + A \cos \alpha (\cos \omega t - 1)}$$



Sol: 1)



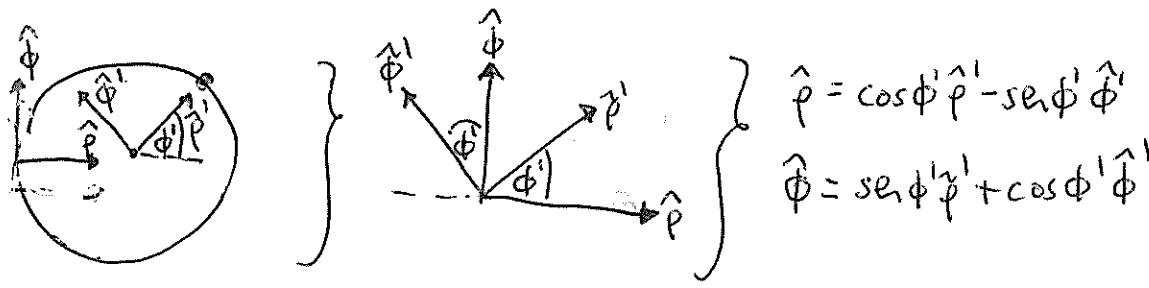
$$\left. \begin{array}{l} \vec{r} = R\hat{p} \\ \vec{v} = R\omega_0\hat{\phi} \\ \vec{a} = -R\omega_0^2\hat{p} \\ \vec{\Omega} = \omega_0\hat{k} \end{array} \right\}$$

2)  $\vec{r}^1 = \hat{p}^1, \hat{\phi}, \hat{k}$ ;  $\vec{S}^1 = \hat{p}^1, \hat{\phi}, \hat{k}$

3) Para el sistema no inercial:

$$\vec{r}^1 = R\hat{p}^1 \quad \vec{v}^1 = R\dot{\phi}\hat{\phi}^1 \quad \vec{a}^1 = -R\ddot{\phi}\hat{p}^1 + R\dot{\phi}^2\hat{\phi}^1$$

4) Relaciono los ejes:

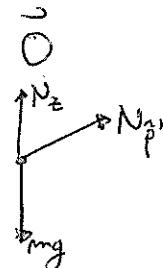


5) Centrifuga:  $m\vec{\Omega} \times (\vec{\Omega} \times \vec{r}) = m\omega_0\hat{k} \times (\omega_0\hat{k} \times R\hat{p}^1) = -mR\omega_0^2\hat{p}^1$

$$\boxed{-\hat{p}^1}$$

Coriolis:  $2m\vec{\Omega} \times \vec{v}^1 = 2m\omega_0\hat{k} \times R\dot{\phi}\hat{\phi}^1 = -2m\omega_0R\hat{p}^1$

Transversal:



6) DCL:  $\vec{F} = N_z\hat{k} - mg\hat{k} + N_p\hat{p}^1$

7)  $m(-R\ddot{\phi}\hat{p}^1 + R\dot{\phi}^2\hat{\phi}^1) = N_z\hat{k} - mg\hat{k} + N_p\hat{p}^1 + R\omega_0^2(\cos\phi\hat{p}^1 - \sin\phi\hat{\phi}^1)m + mR\omega_0^2\hat{p}^1 + 2m\omega_0R\hat{p}^1$

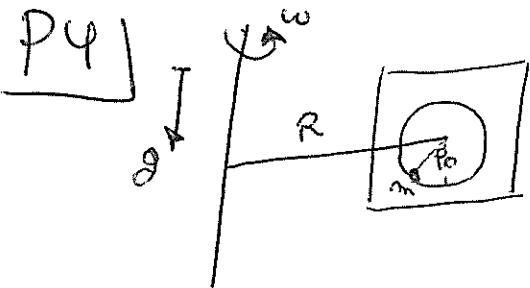
La condición para que llegue al eje sale de la ecuación en  $\hat{\phi}^1$ , al integrar hasta  $\pi$ :

$\hat{\phi}^1: mR\ddot{\phi}^1 = -R\omega_0^2 \sin\phi/m$

$$\int_0^{\pi} \dot{\phi}' d\phi' = -\omega_0^2 \int_0^{\pi} \sin \phi' d\phi'$$

$$-\frac{1}{2} \frac{v_0^2}{R^2} = t U_0^2 \cos \phi' \Big|_0^\pi = \omega_0^2 (-1 - 1) = -2 \omega_0^2$$

$$\Rightarrow \boxed{v_{0,\min}^2 = 4 R^2 \omega_0^2}$$

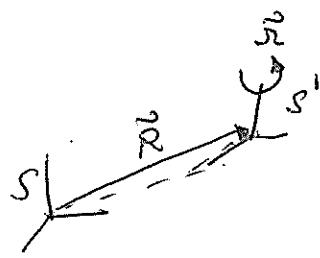


a) Definir  $S, S'$  y calcular pseudo fuerzas

b) Obtener ec. de mov. y las ec. para el ángulo  $\phi = f(\phi)$

c) Bajo qué condiciones  $\phi = 0$  es estable, y calcule freq. de pequeñas oscilaciones

Sol: a) 1)



$$\left. \begin{array}{l} \vec{r} = R \hat{p} \\ \vec{s} = \omega \hat{z} \end{array} \right\} \rightarrow \ddot{\vec{r}} = -R \omega^2 \hat{p}$$

2)  $S: \hat{p}, \hat{\theta}, \hat{z} \quad S': \hat{r}, \hat{\phi}, \hat{z}$

3)  $\vec{r}' = p_0 \hat{r} \quad \vec{v}' = p_0 \dot{\phi} \hat{\phi} \quad \vec{a}' = -p_0 \dot{\phi}^2 \hat{r} + p_0 \ddot{\phi} \hat{\phi}$

4) Ahora relaciona las coordenadas:

$$\left. \begin{array}{l} \hat{p} = \cos \phi \hat{r} + \sin \phi \hat{\theta} \\ \hat{z} = \sin \phi \hat{r} - \cos \phi \hat{\theta} \end{array} \right\} \text{además: } \hat{z}' = \hat{p}$$

5) Pseudo fuerzas:

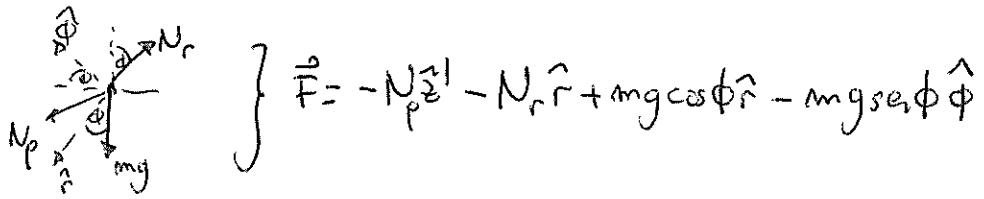
$$\bullet m \ddot{\vec{r}} = -m R \omega^2 \hat{z}'$$

$$\bullet m \vec{r} \times (\vec{r} \times \vec{v}') = m \omega (\sin \phi \hat{\phi} - \cos \phi \hat{r}) \times (m (\sin \phi \hat{\phi} - \cos \phi \hat{r}) \times p_0 \hat{r}) = m \omega^2 p_0 (\sin \phi \hat{\phi} - \cos \phi \hat{r}) \times -\sin \phi \hat{z}' = m \omega^2 p_0 (-\sin^2 \phi \hat{r} - \sin \phi \cos \phi \hat{\phi})$$

$$\bullet 2m \vec{r} \times \vec{v}' = 2m \omega (\sin \phi \hat{\phi} - \cos \phi \hat{r}) \times p_0 \dot{\phi} \hat{\phi} = -2m \omega p_0 \dot{\phi} \cos \phi \hat{z}'$$

$$\bullet m \vec{r} \times \vec{r}' = 0$$

b) 6) DCL:



7) L<sub>2</sub> ec. de mov:

$$\left\{ \begin{array}{l} m(-\rho_0 \dot{\phi}^2 \hat{r} + \rho_0 \ddot{\phi} \hat{\phi}) = -N_p \hat{z} - N_r \hat{r} + mg \cos \phi \hat{r} - mg \sin \phi \hat{\phi} \\ \quad + m R \omega^2 \hat{z} + m \omega_0^2 (\sin^2 \phi \hat{r} + \sin \phi \cos \phi \hat{\phi}) \\ \quad + 2m \omega_0 \dot{\phi} \cos \phi \hat{z} \end{array} \right.$$

Para  $\hat{\phi}$ : 
$$m \rho_0 \ddot{\phi} = -mg \sin \phi + m \omega_0^2 \sin \phi \cos \phi$$

c) Ordenando la ec. anterior:

$$\ddot{\phi} = -\frac{g}{\rho_0} \sin \phi + \omega^2 \sin \phi \cos \phi$$

Para ángulos pequeños:  $\sin \phi \sim \phi$   
 $\cos \phi \sim 1$

$$\Rightarrow \ddot{\phi} = -\frac{g}{\rho_0} \phi + \omega^2 \phi \Rightarrow \ddot{\phi} + \left( \frac{g}{\rho_0} - \omega^2 \right) \phi = 0$$

Para que sea estable:  $\frac{g}{\rho_0} - \omega^2 > 0 \Rightarrow \boxed{\omega_0^2 = \frac{g}{\rho_0} - \omega^2}$   
inestable:  $\frac{g}{\rho_0} - \omega^2 < 0$