

P1 | En cuetres el conmutador $[x_i, L_j]$:

Sol: Dado que $\hat{L} = \hat{r} \times \hat{p}$:

$$\hat{L}_x = \hat{y}\hat{p}_z - \hat{z}\hat{p}_y, \quad \hat{L}_y = \hat{z}\hat{p}_x - \hat{x}\hat{p}_z, \quad \hat{L}_z = \hat{x}\hat{p}_y - \hat{y}\hat{p}_x$$

Entonces:

$$\bullet [\hat{x}, \hat{L}_x] = 0$$

$$\bullet [\hat{z}, \hat{L}_x] = i\hbar\hat{y}$$

$$\bullet [\hat{y}, \hat{L}_y] = \hat{z} \underbrace{[\hat{x}, \hat{p}_x]}_{i\hbar} = i\hbar\hat{z}$$

$$\bullet [\hat{z}, \hat{L}_y] = -i\hbar\hat{x}$$

$$\bullet [\hat{x}, \hat{L}_z] = -\hat{y} [\hat{x}, \hat{p}_x] = -i\hbar\hat{y}$$

$$\bullet [\hat{z}, \hat{L}_z] = 0$$

$$\bullet [\hat{y}, \hat{L}_x] = -i\hbar\hat{z}$$

$$\bullet [\hat{y}, \hat{L}_y] = 0$$

$$\bullet [\hat{y}, \hat{L}_z] = i\hbar\hat{x}$$

De manera general: $[x_i, L_j] = \epsilon_{ijk} i\hbar x_k$

P2 Muestre que:

$$(m-m') \int_0^{2\pi} d\phi \int_0^\pi d\theta \sin\theta Y_{\ell m'}^*(\theta, \phi) \cos\theta Y_{\ell m}(\theta, \phi) = 0$$

y luego que:

$$(m-m')^2 \int_0^{2\pi} d\phi \int_0^\pi d\theta \sin\theta Y_{\ell m'}^*(\theta, \phi) \sin\theta \cos\phi Y_{\ell m}(\theta, \phi) =$$

$$\int_0^{2\pi} d\phi \int_0^\pi d\theta \sin\theta Y_{\ell m'}^*(\theta, \phi) \sin\theta \cos\phi Y_{\ell m}(\theta, \phi)$$

Usado el conmutador $[\hat{L}_z, \hat{x}_i]$

Sol: Primero sacamos elementos de matriz:

$$\bullet \langle \ell' m' | [\hat{L}_z, \hat{x}] | \ell m \rangle = i\hbar \langle \ell' m' | \hat{y} | \ell m \rangle \quad (1a)$$

$$\bullet \langle \ell' m' | [\hat{L}_z, \hat{y}] | \ell m \rangle = -i\hbar \langle \ell' m' | \hat{x} | \ell m \rangle \quad (1b)$$

$$\bullet \langle \ell' m' | [\hat{L}_z, \hat{z}] | \ell m \rangle = 0 \quad (1c)$$

Por otra parte:

$$\bullet \langle \ell' m' | [\hat{L}_z, \hat{x}] | \ell m \rangle = \langle \ell' m' | [\hat{L}_z \hat{x} - \hat{x} \hat{L}_z] | \ell m \rangle =$$

$$= (m' - m) \langle \ell' m' | \hat{x} | \ell m \rangle \quad (2a)$$

$$\bullet \langle \ell' m' | [\hat{L}_z, \hat{y}] | \ell m \rangle = (m' - m) \langle \ell' m' | \hat{y} | \ell m \rangle \quad (2b)$$

$$\bullet \langle \ell' m' | [\hat{L}_z, \hat{z}] | \ell m \rangle = (m' - m) \langle \ell' m' | \hat{z} | \ell m \rangle \quad (2c)$$

Iguando (1c) con (2c):

$$(m' - m) \langle \ell' m' | \hat{z} | \ell m \rangle = 0$$

$$(m' - m) \int_0^{2\pi} d\phi \int_0^\pi d\theta \sin\theta \underbrace{\langle \ell' m' | z | \ell m \rangle}_{Y_{\ell m'}^* \cos\theta} \underbrace{\langle \theta, \phi | \ell m \rangle}_{Y_{\ell m}} = 0$$

de manera análoga:

$$(m'-m) \langle l'm' | \hat{x} | l m \rangle = i \hbar \langle l'm' | \hat{y} | l m \rangle$$

$$(m'-m) \langle l'm' | \hat{y} | l m \rangle = -i \hbar \langle l'm' | \hat{x} | l m \rangle$$

justando ambos:

$$(m'-m) \underbrace{(m'-m) \langle l'm' | \hat{x} | l m \rangle}_{\langle l'm' | \hat{y} | l m \rangle} / i \hbar = -i \hbar \langle l'm' | \hat{x} | l m \rangle$$

$$(m'-m)^2 \langle l'm' | \hat{x} | l m \rangle = \hbar^2 \langle l'm' | \hat{x} | l m \rangle$$

Completado con $\int d\phi \int d\theta |l m \phi \theta\rangle \langle l m \phi \theta|$ se tiene:

$$(m'-m)^2 \int_0^{2\pi} d\phi \int_0^{\pi} d\theta \sin\theta Y_{l'm'}^* \sin\theta \cos\phi Y_{lm} = \int_0^{2\pi} d\phi \int_0^{\pi} d\theta \sin\theta \cos\theta Y_{l'm'}^* Y_{lm}$$

P3 | Se tiene el estado:

$$|\psi\rangle = a|1\ 1\rangle + b|1\ 0\rangle + c|1\ -1\rangle \quad \text{con la notación } |l\ m\rangle$$

donde $|a|^2 + |b|^2 + |c|^2 = 1$

a) Encuentre $\langle L_x \rangle$

b) Encuentre $\langle L^2 \rangle$

c) Encuentre a, b, c para que $\hat{L}_x |\psi\rangle = |\psi\rangle$

Sol: a) Dado que $\hat{L}_+ = \hat{L}_x + i\hat{L}_y \rightarrow \hat{L}_x = \frac{1}{2}\hat{L}_+ + \frac{1}{2}\hat{L}_-$

El valor de expectación $\langle \hat{L}_x \rangle = \frac{1}{2}\langle \hat{L}_+ \rangle + \frac{1}{2}\langle \hat{L}_- \rangle$

• $\langle \hat{L}_+ \rangle = \langle \psi | \hat{L}_+ | \psi \rangle = \langle \psi | [a\sqrt{2-2} |1\ 2\rangle + b\sqrt{2} |1\ 1\rangle + c\sqrt{2} |1\ 0\rangle]$
 $= a^*b\sqrt{2} + b^*c\sqrt{2} = \sqrt{2}(a^*b + b^*c)$

• $\langle \hat{L}_- \rangle = \langle \psi | \hat{L}_- | \psi \rangle = \langle \psi | [a\sqrt{2} |1\ 0\rangle + b\sqrt{2} |1\ -1\rangle + c\sqrt{2-2} |1\ -2\rangle]$
 $= b^*a\sqrt{2} + c^*b\sqrt{2} = \sqrt{2}(b^*a + c^*b)$

Entonces:

$$\langle \hat{L}_x \rangle = \frac{\sqrt{2}}{2} \left(\underbrace{a^*b + b^*a}_{2\text{Re}(a^*b)} + \underbrace{b^*c + c^*b}_{2\text{Re}(b^*c)} \right)$$

$$\boxed{\langle \hat{L}_x \rangle = \sqrt{2}(\text{Re}(a^*b) + \text{Re}(b^*c))}$$

b) Es más directo verlo:

$$\langle \hat{L}^2 \rangle = \langle \psi | \hat{L}^2 | \psi \rangle = \langle \psi | [2a |1\ 1\rangle + 2b |1\ 0\rangle + 2c |1\ -1\rangle] =$$
$$= 2 \underbrace{(a^*a + b^*b + c^*c)}_1 \Rightarrow \boxed{\langle \hat{L}^2 \rangle = 2}$$

$$c) \quad \hat{L}_x |\psi\rangle = |\psi\rangle$$
$$\frac{1}{2} (\hat{L}_+ + \hat{L}_-) |\psi\rangle = |\psi\rangle$$

$$\frac{1}{2} [b\sqrt{2}|1\ 1\rangle + c\sqrt{2}|1\ 0\rangle + a\sqrt{2}|1\ 0\rangle + b\sqrt{2}|1\ -1\rangle] = a|1\ 1\rangle + b|1\ 0\rangle + c|1\ -1\rangle$$

$$\rightarrow \frac{b\sqrt{2}}{2} = a \Rightarrow b = \sqrt{2}a$$

$$\rightarrow \frac{b\sqrt{2}}{2} = c \Rightarrow b = \sqrt{2}c$$

Normalizaremos:

$$|a|^2 + |b|^2 + |c|^2 = b^2 \left[\frac{1}{2} + 1 + \frac{1}{2} \right] = 2b^2 = 1 \Rightarrow \boxed{b = \frac{1}{\sqrt{2}}}$$

Entonces: $\boxed{a = \frac{1}{2}}$, $\boxed{c = \frac{1}{2}}$

P4 La función de onda de una partícula sujeta a un potencial central $V(r)$ es:

$$\Psi(\vec{x}) = (x+y+3z)f(r)$$

- a) Ψ es autofunción de \hat{L}^2 . Si lo es, encuentre l .
 Si no lo es, ¿qué valores posibles de l se podrían obtener si se mide \hat{L}^2 ?
- b) Encuentre la probabilidad de encontrar los estados m posibles.

Sol: Se tiene que:

$$\hat{L}^2 = \left[\frac{1}{\sin^2\theta} \frac{\partial^2}{\partial \phi^2} + \frac{1}{\sin\theta} \frac{\partial}{\partial \theta} \left(\sin\theta \frac{\partial}{\partial \theta} \right) \right]$$

Escribiremos Ψ en esféricas:

$$\Psi(\vec{x}) = r (\cos\phi \sin\theta + \sin\phi \sin\theta + 3\cos\theta) f(r)$$

Como queremos estudiar $\hat{L}^2(\Psi) \rightarrow \langle \vec{x} | \hat{L}^2(\Psi) = \hat{L}^2 \Psi(\vec{x})$

$$\bullet \frac{1}{\sin^2\theta} \frac{\partial^2 \Psi}{\partial \phi^2} = r f(r) \frac{\partial}{\partial \phi} [\cos\phi - \sin\phi] \frac{\sin\theta}{\sin^2\theta} = \frac{r f(r)}{\sin\theta} [-\sin\phi - \cos\phi]$$

$$= \frac{r f(r)}{\sin\theta} (\sin\phi + \cos\phi)$$

$$\bullet \frac{1}{\sin\theta} \frac{\partial}{\partial \theta} \left(\sin\theta \frac{\partial \Psi}{\partial \theta} \right) = \frac{1}{\sin\theta} \frac{\partial}{\partial \theta} [\sin\theta (\cos\phi \cos\theta + \sin\phi \cos\theta - 3\sin\theta)] r f(r) =$$

$$= \frac{r f(r)}{\sin\theta} \left[(\cos\phi + \sin\phi) \underbrace{(\cos^2\theta - \sin^2\theta)}_{1 - 2\sin^2\theta} - 6 \sin\theta \cos\theta \right]$$

entonces:

$$L^2 \psi = \frac{r f(r)}{\sin \theta} \left[(\cancel{\sin \phi + \cos \phi}) - (\cancel{\sin \phi + \cos \phi}) (1 - 2 \sin^2 \theta) + 6 \sin \theta \cos \theta \right] =$$

$$= r f(r) [6 \cos \theta + 2 \sin \theta (\cos \phi + \sin \phi)] =$$

$$= 2r [\cos \phi \sin \theta + \sin \phi \sin \theta + 3 \cos \theta] f(r) = 2\psi$$

Por tanto sí es auto estado. Además:

$$l(l+1) = 2 \Rightarrow \boxed{l=1}$$

b) Dado que $l=1$, se tiene $m=-1, 0, 1$.

Vamos escribir $\psi(\vec{r})$ como función de esféricos armónicos:

$$Y_{1,-1} = \frac{1}{2} \sqrt{\frac{3}{2\pi}} e^{-i\phi} \sin \theta = \frac{1}{2} \sqrt{\frac{3}{2\pi}} \frac{(x-iy)}{r}$$

$$Y_{1,0} = \frac{1}{2} \sqrt{\frac{3}{\pi}} \cos \theta = \frac{1}{2} \sqrt{\frac{3}{\pi}} \frac{z}{r}$$

$$Y_{1,1} = -\frac{1}{2} \sqrt{\frac{3}{2\pi}} e^{i\phi} \sin \theta = -\frac{1}{2} \sqrt{\frac{3}{2\pi}} \frac{(x+iy)}{r}$$

de donde sacamos:

$$z = r \sqrt{\frac{4\pi}{3}} Y_{1,0}$$

$$x = r \sqrt{\frac{2\pi}{3}} (Y_{1,-1} - Y_{1,1})$$

$$y = ir \sqrt{\frac{2\pi}{3}} (Y_{1,-1} + Y_{1,1})$$

Entonces:

$$\psi = \sqrt{\frac{2\pi}{3}} r f(r) [3\sqrt{2} Y_{1,0} + (1+i) Y_{1,-1} + (i-1) Y_{1,1}]$$

dado que nos interesa la parte de m :

$$|\psi_m\rangle = A [3\sqrt{2}|0\rangle + (1+i)|-1\rangle + (i-1)|+1\rangle]$$

Normalizando:

$$1 = A^2 (18 + 2 + 2) = A^2 \cdot 22 \Rightarrow A = \frac{1}{\sqrt{22}}$$

por tanto:

$$|\psi_m\rangle = \frac{3}{\sqrt{11}}|0\rangle + \frac{(1+i)}{\sqrt{22}}|-1\rangle + \frac{(i-1)}{\sqrt{22}}|+1\rangle$$

luego las probabilidades que este en cada estado de m :

$$P_0 = \frac{9}{11}$$

$$P_{+1} = \frac{1}{11}$$

$$P_{-1} = \frac{1}{11}$$