

P1) Evolúe $[x(t), x(0)]$ en el cuadro de Heisenberg para una partícula libre.

Sol: Para una partícula libre: $\hat{H} = \frac{\hat{p}^2}{2m}$

Las ecuaciones de movimiento:

$$\frac{d\hat{p}}{dt} = \frac{1}{i\hbar} [\hat{p}, \hat{H}] = -\frac{\partial \hat{H}}{\partial \hat{x}}$$

$$\frac{d\hat{x}}{dt} = \frac{1}{i\hbar} [\hat{x}, \hat{H}] = \frac{\partial \hat{H}}{\partial \hat{p}}$$

$$\rightarrow \frac{d\hat{p}}{dt} = 0 \Rightarrow p(t) = p(0)$$

$$\rightarrow \frac{d\hat{x}}{dt} = \frac{\hat{p}(t)}{m} = \frac{p(0)}{m} \Rightarrow x(t) = \frac{p(0)}{m}t + x(0)$$

entonces:

$$[x(t), x(0)] = \left[\frac{p(0)}{m}t + x(0), x(0) \right] = \frac{t}{m} [p(0), x(0)] = \boxed{-\frac{i\hbar t}{m}}$$

P2) Considere el hamiltoniano en 3D:

$$\hat{H} = \frac{\hat{p}^2}{2m} + V(x)$$

calculando $[\vec{x} \cdot \vec{p}, H]$, obtenga:

$$\frac{d}{dt} \langle \vec{x} \cdot \vec{p} \rangle = \left\langle \frac{p^2}{2m} \right\rangle - \langle \vec{x} \cdot \vec{\nabla} V \rangle$$

Sol: Sa camos el conmutador:

$$\begin{aligned} [\vec{x} \cdot \vec{p}, H] &= \left[\sum_i x_i p_i, H \right] = \sum_i \left(x_i [p_i, H] + [x_i, H] p_i \right) \\ &= \sum_i \left[-i\hbar x_i \frac{\partial V}{\partial x_i} + i\hbar \frac{p_i}{m} p_i \right] = \frac{i\hbar}{m} \vec{p}^2 - i\hbar \vec{x} \cdot \vec{\nabla} V \end{aligned}$$

$$\text{Usamos: } i\hbar \frac{d\hat{A}}{dt} = [\hat{A}, \hat{H}] + i\hbar \frac{\partial \hat{A}}{\partial t}$$

$$\rightarrow \frac{d\langle \vec{x} \cdot \vec{p} \rangle}{dt} = \frac{1}{i\hbar} [\vec{x} \cdot \vec{p}, H] = \frac{\vec{p}^2}{m} - \vec{x} \cdot \vec{\nabla} V$$

$$\Rightarrow \boxed{\frac{d}{dt} \langle \vec{x} \cdot \vec{p} \rangle = \left\langle \frac{p^2}{m} \right\rangle - \langle \vec{x} \cdot \vec{\nabla} V \rangle}$$

P3 Se busca resolver el oscilador armónico en el cuadro de Heisenberg.

a) Encuentre $\hat{X}(t)$ y $\hat{P}(t)$

b) Evoluando el elemento $\langle n | \hat{X}(t) | n \rangle$, encuentre el espectro de energía.

Sol: a) $\hat{H} = \frac{\hat{p}^2}{2m} + \frac{m}{2} \omega_0^2 \hat{x}^2$

De las ecuaciones de movimiento:

(1) $\frac{d\hat{x}}{dt} = \frac{\hat{p}}{m}$, (2) $\frac{d\hat{p}}{dt} = -m\omega_0^2 \hat{x}^2$

$\frac{d^2\hat{x}}{dt^2} = \frac{1}{m} \frac{d\hat{p}}{dt} = -\omega_0^2 \hat{x} \Rightarrow \hat{x}(t) = \hat{A} \cos(\omega_0 t) + \hat{B} \sin(\omega_0 t)$
 $\hat{x}(0) = \hat{x}_0$

derivado \hat{x} :

$\left. \frac{d\hat{x}}{dt} \right|_{t=0} = \frac{\hat{p}(0)}{m} = \frac{\hat{p}_0}{m} = \omega_0 \hat{B} \Rightarrow \hat{B} = \frac{\hat{p}_0}{m\omega_0}$

entonces:

$\hat{x}(t) = \hat{x}_0 \cos(\omega_0 t) + \frac{\hat{p}_0}{m\omega_0} \sin(\omega_0 t)$ (3)

Usando (1):

$\hat{p}(t) = -m\omega_0 \hat{A} \sin(\omega_0 t) + m\omega_0 \hat{B} \cos(\omega_0 t)$

$\hat{p}(t) = -m\omega_0 \hat{x}_0 \sin(\omega_0 t) + \hat{p}_0 \cos(\omega_0 t)$ (4)

b) el elemento de matriz $\langle n' | \hat{X}(t) | n \rangle$:

$$\langle n' | \hat{X}(t) | n \rangle = \langle n' | e^{\frac{i\hat{H}t}{\hbar}} \hat{X}(0) e^{-\frac{i\hat{H}t}{\hbar}} | n \rangle = e^{i(E_{n'} - E_n)t/\hbar} \langle n' | \hat{X}_0 | n \rangle$$

usando (3):

$$\begin{aligned} \langle n' | \hat{X}(t) | n \rangle &= \langle n' | \hat{X}_0 \cos(\omega_0 t) | n \rangle + \langle n' | \frac{\hat{P}_0}{m\omega_0} \sin(\omega_0 t) | n \rangle = \\ &= \frac{1}{2} \left[e^{i\omega_0 t} (\langle n' | \hat{X}_0 | n \rangle - i \langle n' | \frac{\hat{P}_0}{m\omega_0} | n \rangle) \right. \\ &\quad \left. + e^{-i\omega_0 t} (\langle n' | \hat{X}_0 | n \rangle + i \langle n' | \frac{\hat{P}_0}{m\omega_0} | n \rangle) \right] \end{aligned}$$

Si $E_{n'} > E_n$:

$$* \langle n' | \hat{X}_0 | n \rangle + i \langle n' | \frac{\hat{P}_0}{m\omega_0} | n \rangle = 0$$

$$* e^{i(E_{n'} - E_n)t/\hbar} = e^{i\omega_0 t}$$

entonces: $E_{n'} = E_n + \hbar\omega_0$

Mientras $\langle n' | \hat{X}_0 | n \rangle \neq 0$ se puede seguir la secuencia para energías mayores indefinidamente.

Si $E_{n'} < E_n$: $* \langle n' | \hat{X}_0 | n \rangle - i \langle n' | \frac{\hat{P}_0}{m\omega_0} | n \rangle = 0$

$$* e^{i(E_{n'} - E_n)t/\hbar} = e^{-i\omega_0 t}$$

entonces: $E_{n'} = E_n - \hbar\omega_0$

No se puede bajar indefinidamente ya que las energías deben ser positivas.

Denotando al estado fundamental como $|0\rangle$:

$$E_n = E_0 + n\hbar\omega_0$$

Al examinar el caso $\langle n | \hat{X}(t) | n \rangle$, se ve que debe ser cero porque \hat{x}_0 no tiene dependencia temporal. Por tanto:

$$\langle n | \underbrace{\left(\hat{x}_0 + i \frac{\hat{p}_0}{m\omega_0} \right)}_0 | 0 \rangle = 0$$

$$\Rightarrow \left(\hat{x}_0 - i \frac{\hat{p}_0}{m\omega_0} \right) \left(\hat{x}_0 + i \frac{\hat{p}_0}{m\omega_0} \right) | 0 \rangle = 0$$

$$\left[\hat{x}_0^2 + \frac{\hat{p}_0^2}{m^2\omega_0^2} + \hat{x}_0 \frac{i\hat{p}_0}{m\omega_0} - \frac{i\hat{p}_0}{m\omega_0} \hat{x}_0 \right] | 0 \rangle = 0$$

$$\left[\hat{x}_0^2 + \frac{\hat{p}_0^2}{m^2\omega_0^2} + \frac{i}{m\omega_0} \underbrace{[\hat{x}_0, \hat{p}_0]}_{i\hbar} \right] | 0 \rangle = 0 \quad / \cdot \frac{1}{2} m\omega_0^2$$

$$\left[\frac{\hat{p}_0^2}{2m} + \frac{1}{2} m\omega_0^2 \hat{x}_0^2 + \frac{1}{2} i\omega_0 \cdot i\hbar \right] | 0 \rangle = 0$$

$$\left[\hat{H} - \frac{\hbar\omega_0}{2} \right] | 0 \rangle = 0$$

$$E_0 - \frac{\hbar\omega_0}{2} = 0 \quad \Rightarrow \quad E_0 = \frac{\hbar\omega_0}{2}$$

$$\Rightarrow \boxed{E_n = \hbar\omega_0 \left(n + \frac{1}{2} \right)}$$

Pr) Considera el hamiltoniano:

$$\hat{H} = \frac{\hat{p}^2}{2m} + \frac{m}{2}\omega_0^2 \hat{x}^2 + eE\hat{x}$$

Usando $[\hat{p}(t), \hat{x}(t)] = -i\hbar$ encuentre la ec. de movimiento.
Muestre que $\langle x(t_1), x(t_2) \rangle \neq 0$ para $t_1 \neq t_2$

Sol: Tenemos: $\ast \frac{d\hat{x}}{dt} = \frac{\hat{p}}{m}$

$$\ast \frac{d\hat{p}}{dt} = -m\omega_0^2 \hat{x} - eE$$

$$\Rightarrow \frac{d^2\hat{x}}{dt^2} = -\omega_0^2 \hat{x} - \frac{eE}{m} = -\omega_0^2 \left(\hat{x} + \frac{eE}{m\omega_0^2} \right)$$

$$\Rightarrow \hat{x}(t) = \hat{A} \cos(\omega_0 t) + \hat{B} \sin(\omega_0 t) - \frac{eE}{m\omega_0^2}$$

de forma similar a la pregunta anterior:

$$\hat{x}(t) = \left(\hat{x}_0 + \frac{eE}{m\omega_0^2} \right) \cos(\omega_0 t) + \frac{\hat{p}_0}{m\omega_0} \sin(\omega_0 t) - \frac{eE}{m\omega_0^2}$$

Ahora calculamos el conmutador:

$$\begin{aligned} [\hat{x}(t_1), \hat{x}(t_2)] &= \left[\hat{x}_0 \cos(\omega_0 t_1) + \frac{\hat{p}_0}{m\omega_0} \sin(\omega_0 t_1), \hat{x}_0 \cos(\omega_0 t_2) + \frac{\hat{p}_0}{m\omega_0} \sin(\omega_0 t_2) \right] \\ &= \cos(\omega_0 t_1) \sin(\omega_0 t_2) \frac{1}{m\omega_0} [\hat{x}_0, \hat{p}_0] + \sin(\omega_0 t_1) \cos(\omega_0 t_2) \frac{1}{m\omega_0} [\hat{p}_0, \hat{x}_0] \\ &= \frac{i\hbar}{m\omega_0} \sin(\omega_0 (t_2 - t_1)) \end{aligned}$$